

Quantum Channel Simulation and the Channel's Smooth Max-Information (1807.05354)

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Beyond IID 2018, Cambridge



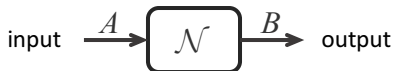
- ⊙ Background
- ⊙ Main results
- ⊙ AEP proof outline
- ⊙ Summary and discussion

Background

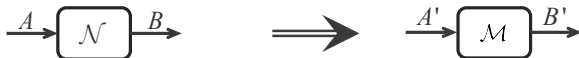
A **quantum channel** is a communication channel which can transmit quantum information. It sends one quantum state to the other.

Mathematically, a quantum channel is characterized by a linear map $\mathcal{N}_{A \rightarrow B}$ that is

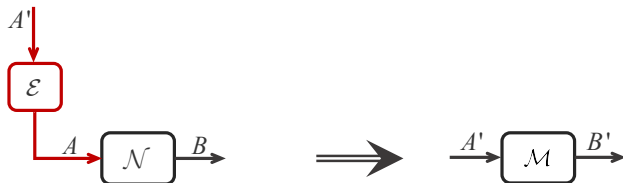
- ⊙ completely positive (**CP**): $\text{id}_k \otimes \mathcal{N}(X_{RA}) \geq 0 \ \forall \ X_{RA} \geq 0$ and $k \in \mathbb{N}$;
- ⊙ trace-preserving (**TP**): $\text{Tr } \mathcal{N}(X) = \text{Tr } X$ for all X .



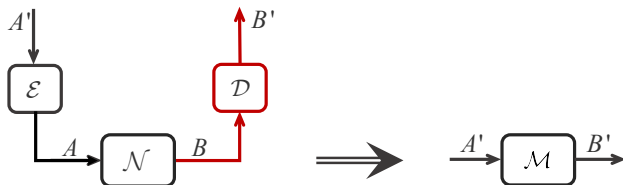
What is channel simulation?



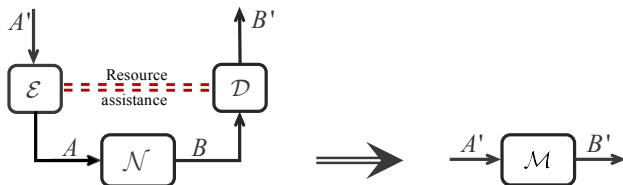
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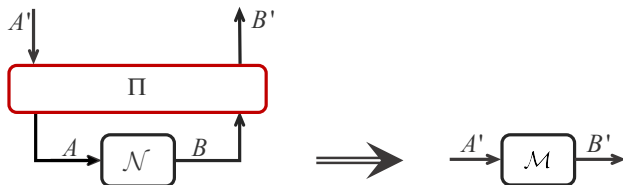
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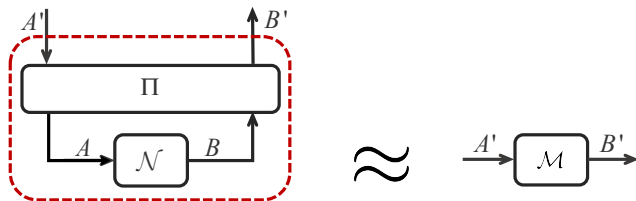
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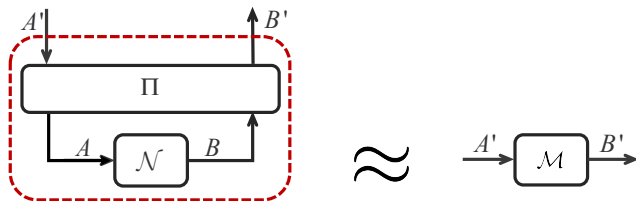


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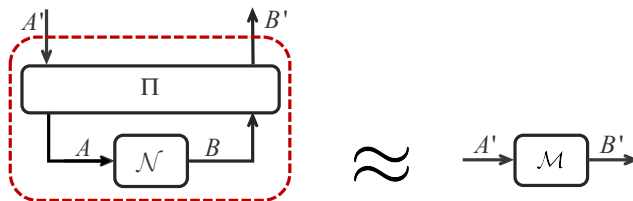


- “Similarity” can be measured via diamond norm [Kitaev, 1997]:

$$\|\mathcal{F}\|_{\diamond} := \sup_k \|\text{id}_k \otimes \mathcal{F}\|_1, \text{ where } \|\cdot\|_1 \text{ induced by the Schatten 1-norm.}$$

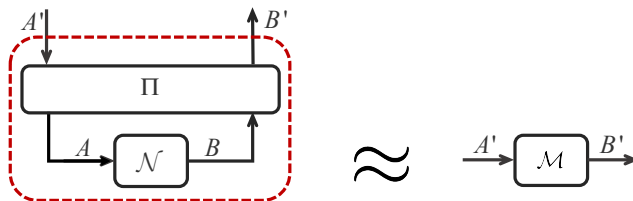
- Operational meaning: minimum error probability p_e to distinguish two quantum channels $\mathcal{N}_1$ and $\mathcal{N}_2$ is given by

$$p_e = \frac{1}{2} \left(1 - \frac{\|\mathcal{N}_1 - \mathcal{N}_2\|_{\diamond}}{2} \right).$$



The minimum error of simulation from $\mathcal{N}$ to $\mathcal{M}$ with Ω -assistance:

$$\omega_{\Omega}(\mathcal{N}, \mathcal{M}) := \frac{1}{2} \inf_{\Pi \in \Omega} \|\Pi \circ \mathcal{N} - \mathcal{M}\|_{\diamond}. \quad (1)$$

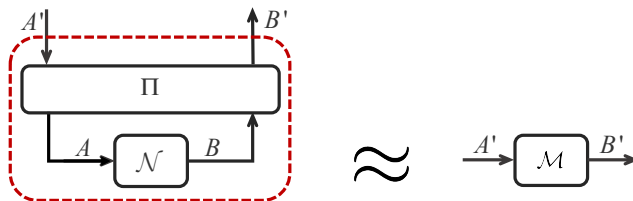


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$$S_{\Omega}(\mathcal{N}, \mathcal{M}) := \liminf_{\varepsilon \rightarrow 0} \left\{ \frac{n}{m} : \omega_{\Omega}(\mathcal{N}^{\otimes n}, \mathcal{M}^{\otimes m}) \leq \varepsilon \right\}. \quad (2)$$



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Different resource assistances can be considered. Here we focus on:

- ⊙ entanglement assistance, $\Omega = \mathcal{E}$;
- ⊙ no-signalling (NS) assistance, $\Omega = \mathcal{NS}$;
- ⊙ $\mathcal{E} \subseteq \mathcal{NS}$.



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$$Q_E(\mathcal{N}) \leq Q_{\text{NS}}(\mathcal{N}) \leq S_{\text{NS}}(\mathcal{N}) \leq S_E(\mathcal{N}). \quad (3)$$



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$$Q_E(\mathcal{N}) \leq Q_{NS}(\mathcal{N}) \leq S_{NS}(\mathcal{N}) \leq S_E(\mathcal{N})$$

$$= \frac{1}{2} I(A : B)_{\mathcal{N}} := \frac{1}{2} \max_{\phi_{AA'}} I(A : B)_{\mathcal{N}_{A' \rightarrow B}(\phi_{AA'})}$$

↑
[Bennett et al., 2002]

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Quantum reverse Shannon theorem (QRTS)

2017 IEEE Information Theory Society Paper Award

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Main results

Most the prior results focus on the **asymptotic** simulation rate.

Our contributions are twofolds:

- ⊙ study the **one-shot** channel simulation task: $\text{id}_m \rightarrow \mathcal{N}$ with NS-assistance;
 - ⊙ introduce a naturally appeared **entropy of a channel** operational meaning + asymptotic equipartition property (AEP)
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Recall the minimum error of simulation:

$$\omega_{\text{NS}}(\text{id}_m, \mathcal{N}) := \frac{1}{2} \inf_{\Pi \in \text{NS}} \|\Pi \circ \text{id}_m - \mathcal{N}\|_{\diamond}. \quad (4)$$

The one-shot quantum simulation cost under NS assistance:

$$S_{\text{NS},\varepsilon}^{(1)}(\mathcal{N}) := \log \min \{m \in \mathbb{N} : \omega_{\text{NS}}(\text{id}_m, \mathcal{N}) \leq \varepsilon\}. \quad (5)$$

Then the asymptotic quantum simulation cost is equivalently given by

$$S_{\text{NS}}(\mathcal{N}) = \lim_{\varepsilon \rightarrow 0} \lim_{n \rightarrow \infty} \frac{1}{n} S_{\text{NS},\varepsilon}^{(1)}(\mathcal{N}^{\otimes n}). \quad (6)$$

The minimum error $\omega_{\text{NS}}(\text{id}_m, \mathcal{N})$ can be given by a SDP,

$$\text{minimize } \lambda \tag{7a}$$

$$\text{subject to } \text{Tr}_B Y_{AB} \leq \lambda \mathbb{1}_A, \tag{7b}$$

$$Y_{AB} \geq J_{\tilde{\mathcal{N}}} - J_{\mathcal{N}}, \quad Y_{AB} \geq 0, \tag{7c}$$

$$J_{\tilde{\mathcal{N}}} \geq 0, \quad \text{Tr}_B J_{\tilde{\mathcal{N}}} = \mathbb{1}_A, \tag{7d}$$

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Sketch of proof: $\omega_{\text{NS}}(\text{id}_m, \mathcal{N}) := \frac{1}{2} \inf_{\Pi \in \text{NS}} \|\Pi \circ \text{id}_m - \mathcal{N}\|_{\diamond},$

⊙ [Leung & Matthews, 2015; Duan & Winter, 2016] $\Pi \in \text{NS}$ iff

$$J_{\Pi} \geq 0, \quad \text{Tr}_{AB'} J_{\Pi} = \mathbb{1}_{A'B},$$

$$\text{Tr}_A J_{\Pi} = \mathbb{1}_{A'}/d_{A'} \otimes \text{Tr}_{AA'} J_{\Pi}, \quad \text{Tr}_{B'} J_{\Pi} = \mathbb{1}_B/d_B \otimes \text{Tr}_{BB'} J_{\Pi}$$

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⊙ [Watrous, 2009] $\frac{1}{2} \|\mathcal{N}_1 - \mathcal{N}_2\|_\diamond$ given by SDP

$$\min\{\lambda : \text{Tr}_B Y_{AB} \leq \lambda \mathbb{1}_A, Y_{AB} \geq J_{\mathcal{N}_1} - J_{\mathcal{N}_2}, Y_{AB} \geq 0\}$$

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⊙ Symmetry of id_m : Choi matrix invariant under $U \otimes \bar{U}$.

The one-shot ε -error quantum simulation cost $S_{\text{NS},\varepsilon}^{(1)}(\mathcal{N})$

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We will need a entropy that looks more similar to $I(A : B)_{\mathcal{N}}$.

The quantum mutual information of a state

$$I(A : B)_\rho := \inf_{\sigma_B} D(\rho_{AB} \| \rho_A \otimes \sigma_B). \quad (10)$$

The max-information of a quantum state [Berta et al., 2011]:

$$I_{\max}(A : B)_\rho := \inf_{\sigma_B} D_{\max}(\rho_{AB} \| \rho_A \otimes \sigma_B), \quad (11)$$

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Other variations can be seen in [Ciganović et al., 2013].

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The smoothed version:

$$I_{\max}^\varepsilon(A : B)_\rho := \inf_{\tilde{\rho} \approx^\varepsilon \rho} I_{\max}(A : B)_{\tilde{\rho}}. \quad (12)$$

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$$I_{\max}^\varepsilon(A : B)_\rho := \inf_{\tilde{\rho} \approx^\varepsilon \rho} I_{\max}(A : B)_{\tilde{\rho}}. \quad (12)$$

The asymptotic equipartition property (AEP) [Berta et al., 2011]:

$$\lim_{\varepsilon \rightarrow 0} \lim_{n \rightarrow \infty} \frac{1}{n} I_{\max}^\varepsilon(A : B)_{\rho^{\otimes n}} = I(A : B)_\rho. \quad (13)$$

The quantum mutual information of a state

$$I(A : B)_\rho := \inf_{\sigma_B} D(\rho_{AB} \| \rho_A \otimes \sigma_B). \quad (10)$$

The max-information of a quantum state [Berta et al., 2011]:

$$I_{\max}(A : B)_\rho := \inf_{\sigma_B} D_{\max}(\rho_{AB} \| \rho_A \otimes \sigma_B), \quad (11)$$

where max-relative entropy [Datta 2009] $D_{\max}(\rho \| \sigma) = \inf\{t | \rho \leq 2^t \sigma\}$.

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- ⊙ We will generalize these notations and results to a **channel's version** and find their connection with the quantum channel simulation task.

Definition

For any quantum channel $\mathcal{N}_{A' \rightarrow B}$ we define the max-information of the channel $\mathcal{N}$ as

$$I_{\max}(A : B)_{\mathcal{N}} := I_{\max}(A : B)_{\mathcal{N}_{A' \rightarrow B}(\Phi_{AA'})}, \quad (14)$$

where $\Phi_{AA'}$ is the maximally entangled state.

We can replace $\Phi_{AA'}$ to any pure state $\phi_{AA'}$ with Schmidt rank $|A'|$.

Definition

The channel's **smooth** max-information is defined by

$$I_{\max}^{\varepsilon}(A : B)_{\mathcal{N}} := \inf_{\substack{\frac{1}{2} \|\tilde{\mathcal{N}} - \mathcal{N}\|_{\diamond} \leq \varepsilon \\ \tilde{\mathcal{N}} \in \text{CTP}(A' : B)}} I_{\max}(A : B)_{\tilde{\mathcal{N}}}, \quad (15)$$

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Monotone under composition with CPTP maps, i.e., for any CPTP maps $\mathcal{N}_{A' \rightarrow B_1}$, $\mathcal{F}_{A'_0 \rightarrow A'_1}$ and $\mathcal{T}_{B_1 \rightarrow B_0}$,

$$I_{\max}^{\varepsilon}(A_0 : B_0)_{\mathcal{T} \circ \mathcal{N} \circ \mathcal{F}} \leq I_{\max}^{\varepsilon}(A_1 : B_1)_{\mathcal{N}}. \quad (16)$$

$$D_{\max}(\mathcal{N}||\mathcal{M}) := D_{\max}(J_{\mathcal{N}}||J_{\mathcal{M}}) \quad (17)$$

$$I_{\max}(A : B)_{\mathcal{N}} = \inf_{\mathcal{M} \in \mathcal{G}} D_{\max}(\mathcal{N}||\mathcal{M}) \quad (18)$$

$\mathcal{G} := \{ \mathcal{M} \in \text{CPTP} : \exists \sigma \text{ s.t. } \mathcal{M}(\rho) = \sigma, \forall \rho \}$ the set of constant channels.

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Max-relative entropy and robustness for states: [Datta, 2009]

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The robustness of a quantum channel $\mathcal{N}$ (see also [Díaz et.al, 2018])

$$\mathcal{R}_g(\mathcal{N}) := \inf \left\{ t \geq 0 \mid \exists \mathcal{M} \in \text{CPTP}(A : B) \text{ s.t. } \frac{\mathcal{N} + t\mathcal{M}}{1+t} \in \mathcal{G} \right\}. \quad (19)$$

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$$I_{\max}(A : B)_{\mathcal{N}} = \log[1 + \mathcal{R}_g(\mathcal{N})]. \quad (20)$$

Theorem

For any quantum channel $\mathcal{N}_{A' \rightarrow B}$ and given error tolerance $\varepsilon \geq 0$, we have

$$S_{\text{NS}, \varepsilon}^{(1)}(\mathcal{N}) = \frac{1}{2} I_{\max}^{\varepsilon}(A : B)_{\mathcal{N}}. \quad (21)$$

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$$\lim_{\varepsilon \rightarrow 0} \lim_{n \rightarrow \infty} \frac{1}{n} I_{\text{max}}^{\varepsilon}(A : B)_{\mathcal{N}^{\otimes n}} = I(A : B)_{\mathcal{N}}. \quad (22)$$

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Sketch of proof:

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \lim_{n \rightarrow \infty} \frac{1}{n} I_{\text{max}}^{\varepsilon}(A : B)_{\mathcal{N}^{\otimes n}} &= 2 \cdot \lim_{\varepsilon \rightarrow 0} \lim_{n \rightarrow \infty} \frac{1}{n} S_{\text{NS},\varepsilon}^{(1)}(\mathcal{N}^{\otimes n}) \\ &= 2 \cdot S_{\text{NS}}(\mathcal{N}) && \text{[by definition]} \\ &= 2 \cdot Q_E(\mathcal{N}) && \text{[Bennett et al., 2014]} \\ &= I(A : B)_{\mathcal{N}} && \text{[Bennett et al., 2002]} \end{aligned}$$

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Note: on the other hand, if we can proof Eq. (22) directly, it implies

$$Q_{\text{E}}(\mathcal{N}) = \underbrace{Q_{\text{NS}}(\mathcal{N}) = S_{\text{NS}}(\mathcal{N})}_{\text{QRST}} \leq S_{\text{E}}(\mathcal{N})$$

AEP proof outline

$$\lim_{\varepsilon \rightarrow 0} \lim_{n \rightarrow \infty} \frac{1}{n} I_{\max}^{\varepsilon}(A : B)_{\mathcal{N}^{\otimes n}} \geq I(A : B)_{\mathcal{N}}$$

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$$I_{\max}^{\varepsilon}(A : B)_{\mathcal{N}^{\otimes n}} = \inf_{\frac{1}{2} \|\tilde{\mathcal{N}}^n - \mathcal{N}^{\otimes n}\|_{\diamond} \leq \varepsilon} I_{\max}(A : B)_{\tilde{\mathcal{N}}_{A' \rightarrow B}^n(\phi_{AA'}^{\otimes n})} \quad [\text{definition}]$$

$$= n I(A : B)_{\mathcal{N}_{A' \rightarrow B}(\phi_{AA'})} \quad [\text{additivity}]$$

$$= n I(A : B)_{\mathcal{N}} \quad [\text{optimal } \phi_{AA'}]$$

$$\lim_{\varepsilon \rightarrow 0} \lim_{n \rightarrow \infty} \frac{1}{n} I_{\max}^{\varepsilon}(A : B)_{\mathcal{N}^{\otimes n}} \geq I(A : B)_{\mathcal{N}}$$

$$\begin{aligned}
 I_{\max}^{\varepsilon}(A : B)_{\mathcal{N}^{\otimes n}} &= \inf_{\frac{1}{2} \|\tilde{\mathcal{N}}^n - \mathcal{N}^{\otimes n}\|_{\diamond} \leq \varepsilon} I_{\max}(A : B)_{\tilde{\mathcal{N}}_{A' \rightarrow B}^n(\phi_{AA'}^{\otimes n})} && [\text{definition}] \\
 &\stackrel{[1]}{\geq} \inf_{\frac{1}{2} \|\tilde{\mathcal{N}}^n - \mathcal{N}^{\otimes n}\|_{\diamond} \leq \varepsilon} I(A : B)_{\tilde{\mathcal{N}}_{A' \rightarrow B}^n(\phi_{AA'}^{\otimes n})} && [D_{\max} \geq D] \\
 &= nI(A : B)_{\mathcal{N}_{A' \rightarrow B}(\phi_{AA'})} && [\text{additivity}] \\
 &= nI(A : B)_{\mathcal{N}} && [\text{optimal } \phi_{AA'}]
 \end{aligned}$$

[1] N. Datta, "Min-and max-relative entropies and a new entanglement monotone", 2009

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 &\stackrel{[1]}{\geq} \inf_{\frac{1}{2} \|\tilde{\mathcal{N}}^n - \mathcal{N}^{\otimes n}\|_{\diamond} \leq \varepsilon} I(A : B)_{\tilde{\mathcal{N}}_{A' \rightarrow B}^n(\phi_{AA'}^{\otimes n})} && [D_{\max} \geq D] \\
 &\stackrel{[2]}{\gtrsim} I(A : B)_{\mathcal{N}_{A' \rightarrow B}^{\otimes n}(\phi_{AA'}^{\otimes n})} && [\text{continuity}] \\
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[1] N. Datta, "Min-and max-relative entropies and a new entanglement monotone", 2009

[2] R. Alicki, M. Fannes, "Continuity of quantum conditional information", 2004

$$\left\| \tilde{\mathcal{N}}^n - \mathcal{N}^{\otimes n} \right\|_{\diamond} = \sup_{\phi_{AA'}^n} \left\| \tilde{\mathcal{N}}_{A' \rightarrow B}^n(\phi_{AA'}^n) - \mathcal{N}_{A' \rightarrow B}^{\otimes n}(\phi_{AA'}^n) \right\|_1 \leq \varepsilon$$

$$\begin{aligned} \left\| \tilde{\mathcal{N}}^n - \mathcal{N}^{\otimes n} \right\|_{\diamond} &= \sup_{\phi_{AA'}^n} \left\| \tilde{\mathcal{N}}_{A' \rightarrow B}^n(\phi_{AA'}^n) - \mathcal{N}_{A' \rightarrow B}^{\otimes n}(\phi_{AA'}^n) \right\|_1 \leq \varepsilon \\ &\quad \uparrow \text{[Christandl, König, Renner 2009]} \\ \left\| \tilde{\mathcal{N}}_{A' \rightarrow B}^n(\omega_{RAA'}^n) - \mathcal{N}_{A' \rightarrow B}^{\otimes n}(\omega_{RAA'}^n) \right\|_1 &\leq \varepsilon (n+1)^{-(|A'|^2-1)} \end{aligned}$$

- ⊙ $\omega_{RAA'}^n$ is a purification of the **de Finetti state**

$$\omega_{AA'}^n := \int \phi_{AA'}^{\otimes n} d(\phi_{AA'}),$$

$d(\cdot)$ measure on normalized pure states induced by Haar measure;

- ⊙ We can make $|R| \leq (n+1)^{|A'|^2-1}$.
(see e.g. [Berta, Christandl, Renner 2011])

$$\begin{aligned} \|\tilde{\mathcal{N}}^n - \mathcal{N}^{\otimes n}\|_{\diamond} &= \sup_{\phi_{AA'}^n} \left\| \tilde{\mathcal{N}}_{A' \rightarrow B}^n(\phi_{AA'}^n) - \mathcal{N}_{A' \rightarrow B}^{\otimes n}(\phi_{AA'}^n) \right\|_1 \leq \varepsilon \\ &\quad \Uparrow \text{ [Christandl, König, Renner 2009]} \\ \left\| \tilde{\mathcal{N}}_{A' \rightarrow B}^n(\omega_{RAA'}^n) - \mathcal{N}_{A' \rightarrow B}^{\otimes n}(\omega_{RAA'}^n) \right\|_1 &\leq \varepsilon(n+1)^{-(|A'|^2-1)} \end{aligned}$$

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We are ready to prove

$$\lim_{\varepsilon \rightarrow 0} \lim_{n \rightarrow \infty} \frac{1}{n} I_{\max}^{\varepsilon}(A : B)_{\mathcal{N}^{\otimes n}} \leq I(A : B)_{\mathcal{N}}$$

$$I_{\max}^{\varepsilon}(A : B)_{\mathcal{N}^{\otimes n}} = \inf_{\frac{1}{2} \|\tilde{\mathcal{N}}^n - \mathcal{N}^{\otimes n}\|_{\diamond} \leq \varepsilon} I_{\max}(AR : B)_{\tilde{\mathcal{N}}_{A' \rightarrow B}^n(\omega_{A'AR}^n)} \quad [\text{definition}]$$

$$\begin{aligned}
 I_{\max}^{\varepsilon}(A : B)_{\mathcal{N}^{\otimes n}} &= \inf_{\frac{1}{2} \|\tilde{\mathcal{N}}^n - \mathcal{N}^{\otimes n}\|_{\diamond} \leq \varepsilon} I_{\max}(AR : B)_{\tilde{\mathcal{N}}_{A' \rightarrow B}^n(\omega_{A'AR}^n)} && [\text{definition}] \\
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 &= \inf_{\substack{\frac{1}{2} \|\sigma_{BAR}^n - \mathcal{N}^{\otimes n}(\omega_{A'AR}^n)\|_1 \leq \varepsilon \\ \sigma_{AR}^n = \omega_{AR}^n}} I_{\max}(AR : B)_{\sigma_{BAR}^n} && [\text{partial smooth}]
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- [1] A. Anshu, M. Berta, R. Jain, and M. Tomamichel, arXiv:1807.05630. Talk on Thursday by Mario.
- [2] M. Berta, M. Christandl, and R. Renner, "The quantum reverse Shannon theorem based on one-shot information theory", 2011.

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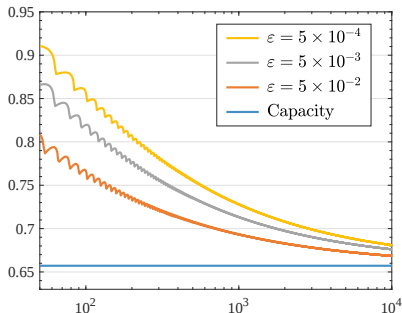
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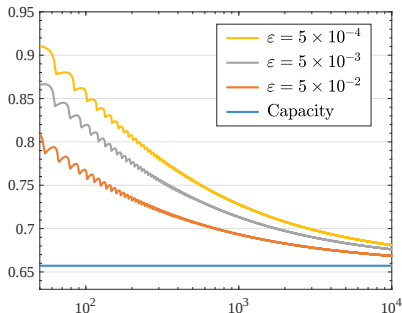
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Depolarizing channel: $\mathcal{N}(\rho) = (1 - p)\rho + p \cdot \mathbb{1}/d$, $p = 0.15$.



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 S_{NS, \epsilon}^{(1)}(\mathcal{N}) &= \frac{1}{2} I_{\max}^{\epsilon}(A : B)_{\mathcal{N}} \\
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 &\rightarrow \frac{1}{2} I(A : B)_{\mathcal{N}}.
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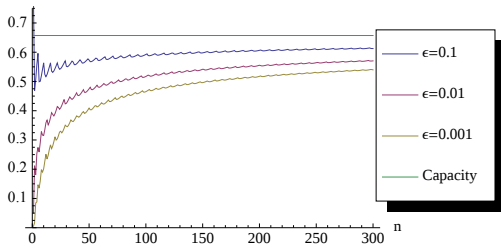
Depolarizing channel: $\mathcal{N}(\rho) = (1 - p)\rho + p \cdot \mathbb{1}/d$, $p = 0.15$.



$$\begin{aligned} S_{NS, \epsilon}^{(1)}(\mathcal{N}) &= \frac{1}{2} I_{\max}^{\epsilon}(A : B)_{\mathcal{N}} \\ &= \frac{1}{2} \min_{\mathcal{M} \in \mathcal{G}} D_{\max}^{\epsilon}(\mathcal{N} \| \mathcal{M}) \\ &\rightarrow \frac{1}{2} I(A : B)_{\mathcal{N}}. \end{aligned}$$

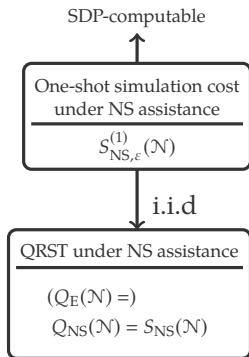
Matthews-Werner
converse bound 2014

$$\begin{aligned} MW(\mathcal{N}, \epsilon) &= \\ &\frac{1}{2} \min_{\mathcal{M} \in \mathcal{G}} D_H^{\epsilon}(\mathcal{N} \| \mathcal{M}) \\ &\rightarrow \frac{1}{2} I(A : B)_{\mathcal{N}} \end{aligned}$$

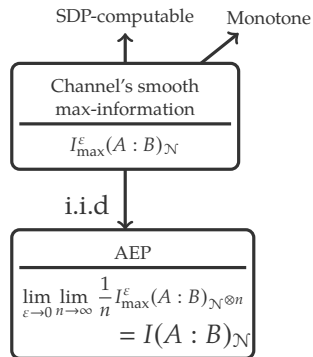




Channel simulation task



Channel's max-information



- ⊙ How to better characterize the one-shot entanglement-assisted channel simulation cost?

- ⊙ How to better characterize the one-shot entanglement-assisted channel simulation cost?
- ⊙ More general form of channel AEP?
(presented by Andreas in JILA workshop open problem session)
Channel divergence:

$$D(\mathcal{N}||\mathcal{M}) := \max_{\phi_{AA'}} D(\mathcal{N}_{A' \rightarrow B}(\phi_{AA'}) || \mathcal{M}_{A' \rightarrow B}(\phi_{AA'})). \quad (23)$$

What we already know for $\mathcal{G}_n$ the set of constant channels:

$$\min_{\mathcal{M}^n \in \mathcal{G}_n} D_{\max}^{\varepsilon}(\mathcal{N}^{\otimes n} || \mathcal{M}^n) \xrightarrow{\text{this work}} \min_{\mathcal{M} \in \mathcal{G}} D(\mathcal{N} || \mathcal{M}) \quad \checkmark \quad (24)$$

$$\min_{\mathcal{M}^n \in \mathcal{G}_n} D_H^{\varepsilon}(\mathcal{N}^{\otimes n} || \mathcal{M}^n) \xrightarrow{\text{MW bound}} \min_{\mathcal{M} \in \mathcal{G}} D(\mathcal{N} || \mathcal{M}) \quad \checkmark \quad (25)$$

What if $\mathcal{G}_n = \{\mathcal{M}^{\otimes n}\}$? Quantum channel Stein's lemma?

Thanks for your attention!

See arXiv:1807.05354

for more details