## Quantum Channel Simulation and the Channel's Smooth Max-Information

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In quantum information theory, a quantum channel is a communication channel which can transmit quantum information. It sends one quantum state to the other.

Mathematically, a quantum channel is characterized by a linear map  $\mathcal{N}_{A \rightarrow B}$  that is

- ◎ completely positive **(CP)**:  $id_k \otimes \mathcal{N}(X_{RA}) \ge 0$  for all  $X_{RA} \ge 0$  and  $k \in \mathbb{N}$ ;
- ◎ trace-preserving **(TP)**:  $\operatorname{Tr} \mathcal{N}(X) = \operatorname{Tr} X$  for all *X*.

input 
$$\xrightarrow{A}$$
  $\mathcal{N}$   $\xrightarrow{B}$  output















• "Similarity" can be measured via diamond norm [Kitaev, 1997]:

 $\|\mathcal{F}\|_{\diamond} \coloneqq \sup_{k} \|\mathrm{id}_k \otimes \mathcal{F}\|_1, \quad \|\cdot\|_1 \text{ induced by the Schatten 1-norm.}$ 

<sup>(3)</sup> Nice operational meaning: minimum error probability  $p_e$  to distinguish two quantum channels  $N_1$  and  $N_2$  is given by

$$p_e = \frac{1}{2} \left( 1 - \frac{\|\mathcal{N}_1 - \mathcal{N}_2\|_\diamond}{2} \right).$$



The minimum error of simulation from  ${\mathbb N}$  to  ${\mathbb M}$  with  $\Omega\text{-assistance}$  is defined as

$$\omega_{\Omega}(\mathbb{N}, \mathbb{M}) \coloneqq \frac{1}{2} \inf_{\Pi \in \Omega} \|\Pi \circ \mathbb{N} - \mathbb{M}\|_{\diamond}.$$
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The channel simulation rate from  ${\mathcal N}$  to  ${\mathcal M}$  with  $\Omega\text{-assistance}$  is defined as

$$S_{\Omega}(\mathbb{N}, \mathbb{M}) := \lim_{\varepsilon \to 0} \inf \left\{ \frac{n}{m} : \omega_{\Omega}\left(\mathbb{N}^{\otimes n}, \mathbb{M}^{\otimes m}\right) \le \varepsilon \right\}.$$
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Different resource assistances can be considered. Here we focus on:

- entanglement assistance,  $\Omega = E$ ;
- ◎ no-signalling (NS) assistance, Ω = NS;
- $\odot$  E  $\subseteq$  NS.



**Q**: What is the optimal rate to simulate the identity channel id<sub>2</sub> via given channel N? In the framework of channel simulation, we have  $Q_{\Omega}(N) = S_{\Omega}(N, id_2)^{-1}$ .



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Operationally, it holds

$$Q_{\rm E}(\mathcal{N}) \le Q_{\rm NS}(\mathcal{N}) \le S_{\rm NS}(\mathcal{N}) \le S_{\rm E}(\mathcal{N}).$$
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### $Q_{\mathrm{E}}(\mathcal{N}) \leq Q_{\mathrm{NS}}(\mathcal{N}) \leq S_{\mathrm{NS}}(\mathcal{N}) \leq S_{\mathrm{E}}(\mathcal{N})$

$$\int_{[\text{Bennett et al., 2002}]} \frac{1}{2} \prod_{\rho_{A'}} I(A:B)_{\mathcal{N}_{A' \to B}}(\phi_{AA'})$$

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Quantum reverse Shannon theorem (QRTS)

2017 IEEE Information Theory Society Paper Award

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- ◎ study the **one-shot** channel simulation scenario, especially  $id_m \rightarrow N$  with NS-assistance;
- ◎ introduce a naturally arisen **entropy of a channel**.

Recall the minimum error of simulation:

$$\omega_{\rm NS}\left({\rm id}_m, \mathcal{N}\right) := \frac{1}{2} \inf_{\Pi \in {\rm NS}} \|\Pi \circ {\rm id}_m - \mathcal{N}\|_{\diamond}. \tag{4}$$

The one-shot quantum simulation cost under NS assistance is defined as

$$S_{\mathrm{NS},\varepsilon}^{(1)}(\mathcal{N}) := \log \min \left\{ m \in \mathbb{N} : \omega_{\mathrm{NS}}(\mathrm{id}_m, \mathcal{N}) \le \varepsilon \right\}.$$
(5)

Then the asymptotic quantum simulation cost is equivalently given by

$$S_{\rm NS}(\mathcal{N}) = \lim_{\varepsilon \to 0} \lim_{n \to \infty} \frac{1}{n} S_{\rm NS,\varepsilon}^{(1)} \left( \mathcal{N}^{\otimes n} \right).$$
(6)

$$\begin{array}{ll} \text{minimize} & \lambda & (7a) \\ \text{subject to} & \operatorname{Tr}_{B'} Y_{A'B'} \leq \lambda \mathbbm{1}_{A'}, & (7b) \\ & Y_{A'B'} \geq J_{\widetilde{\mathcal{N}}} - J_{\mathcal{N}}, Y_{A'B'} \geq 0, & (7c) \\ & J_{\widetilde{\mathcal{N}}} \geq 0, \operatorname{Tr}_{B'} J_{\widetilde{\mathcal{N}}} = \mathbbm{1}_{A'}, & (7d) \\ & J_{\widetilde{\mathcal{N}}} \leq \mathbbm{1}_{A'} \otimes V_{B'}, \operatorname{Tr} V_{B'} = m^2. & (7e) \end{array}$$

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Sketch of proof: In the definition  $\omega_{NS}(id_m, N) := \frac{1}{2} \inf_{\Pi \in NS} ||\Pi \circ id_m - N||_{\diamond}$ , note that

◎  $\Pi \in \text{NS}$  if and only if [Leung and Matthews, 2015; Duan and Winter, 2016]  $\underline{J_{\Pi} \ge 0, \text{ Tr}_{AB'} J_{\Pi} = \mathbb{1}_{A'B}}_{I_{\Pi}}; \text{ Tr}_{A} J_{\Pi} = \frac{\mathbb{1}_{A'}}{d_{A'}} \otimes \text{Tr}_{AA'} J_{\Pi}; \text{ Tr}_{B'} J_{\Pi} = \frac{\mathbb{1}_{B}}{d_{B}} \otimes \text{Tr}_{BB'} J_{\Pi}.$ 

$$\begin{array}{ll} \text{minimize } \lambda & (7a) \\ \text{subject to } & \operatorname{Tr}_{B'} Y_{A'B'} \leq \lambda \mathbbm{1}_{A'}, & (7b) \\ & Y_{A'B'} \geq J_{\widetilde{N}} - J_{\mathcal{N}}, Y_{A'B'} \geq 0, & (7c) \\ & J_{\widetilde{N}} \geq 0, & \operatorname{Tr}_{B'} J_{\widetilde{N}} = \mathbbm{1}_{A'}, & (7d) \\ & J_{\widetilde{N}} \leq \mathbbm{1}_{A'} \otimes V_{B'}, & \operatorname{Tr} V_{B'} = m^2. & (7e) \end{array}$$

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$$\label{eq:main_states} @ \ \frac{1}{2} \|\mathcal{N}_1 - \mathcal{N}_2\|_{\diamond} = \min \left\{ \lambda : \mathrm{Tr}_B \; Y_{AB} \leq \lambda \mathbb{1}_A, Y_{AB} \geq J_{\mathcal{N}_1} - J_{\mathcal{N}_2}, Y_{AB} \geq 0 \right\} [ \mathrm{Watrous}, 2009].$$

$$\begin{array}{ll} \text{minimize} & \lambda & (7a) \\ \text{subject to} & \operatorname{Tr}_{B'} Y_{A'B'} \leq \lambda \mathbbm{1}_{A'}, & (7b) \\ & Y_{A'B'} \geq J_{\widetilde{\mathcal{N}}} - J_{\mathcal{N}}, Y_{A'B'} \geq 0, & (7c) \\ & J_{\widetilde{\mathcal{N}}} \geq 0, \operatorname{Tr}_{B'} J_{\widetilde{\mathcal{N}}} = \mathbbm{1}_{A'}, & (7d) \\ & J_{\widetilde{\mathcal{N}}} \leq \mathbbm{1}_{A'} \otimes V_{B'}, \operatorname{Tr} V_{B'} = m^2. & (7e) \end{array}$$

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- $\label{eq:constraint} \textcircled{0} \quad \frac{1}{2} \|\mathcal{N}_1 \mathcal{N}_2\|_{\diamond} = \min\left\{\lambda: \mathrm{Tr}_B \; Y_{AB} \leq \lambda \mathbbm{1}_A, Y_{AB} \geq J_{\mathcal{N}_1} J_{\mathcal{N}_2}, Y_{AB} \geq 0\right\} \; [\mathrm{Watrous}, 2009].$

◎ Symmetry of  $id_m$ : its Choi matrix is invariant under  $U \otimes \overline{U}$  for any unitary U.

The one-shot  $\varepsilon$ -error quantum simulation cost  $S_{NS,\varepsilon}^{(1)}$  ( $\mathbb{N}$ ) is given by the following SDP,

$$\frac{1}{2}\log$$
 minimize Tr  $V_{B'}$  (8a)

subject to 
$$\operatorname{Tr}_{B'} Y_{A'B'} \leq \varepsilon \mathbb{1}_{A'}$$
, (8b)

$$Y_{A'B'} \ge J_{\widetilde{\mathcal{N}}} - J_{\mathcal{N}}, \ Y_{A'B'} \ge 0, \tag{8c}$$

$$J_{\widetilde{N}} \ge 0, \ \operatorname{Tr}_{B'} J_{\widetilde{N}} = \mathbb{1}_{A'},$$
(8d)

$$J_{\widetilde{\mathcal{N}}} \le \mathbb{1}_{A'} \otimes V_{B'}. \tag{8e}$$

It reduces to the zero-error case when  $\varepsilon = 0$  [Duan and Winter, 2016],

$$S_{\text{NS},0}^{(1)}(\mathbb{N}) = \frac{1}{2} \log \min \left\{ \text{Tr} \, V_{B'} : J_{\mathcal{N}} \le \mathbb{1}_{A'} \otimes V_{B'} \right\} =: \frac{1}{2} H_{\min} \left( A | B \right)_{J_{\mathcal{N}}}.$$
(9)

Since the zero-error cost is additive, we have

-1

$$S_{\rm NS,0}(\mathbb{N}) := \lim_{n \to \infty} \frac{1}{n} S_{\rm NS,0}^{(1)}(\mathbb{N}^{\otimes n}) = S_{\rm NS,0}^{(1)}(\mathbb{N}).$$
(10)

The max-information of a quantum state [Berta et al., 2011]:

$$I_{\max}(A:B)_{\rho} := \inf_{\sigma_B} D_{\max}\left(\rho_{AB} \| \rho_A \otimes \sigma_B\right), \tag{11}$$

where the max-relative entropy [Datta, 2009]  $D_{\max}(\rho \| \sigma) := \log \inf\{t \mid \rho \le t \cdot \sigma\}$ . The smoothed version:

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The asymptotic equipartition property (AEP) holds [Berta et al., 2011]:

$$\lim_{\varepsilon \to 0} \lim_{n \to \infty} \frac{1}{n} I^{\varepsilon}_{\max} \left( A : B \right)_{\rho^{\otimes n}} = I \left( A : B \right)_{\rho}, \tag{13}$$

where the quantum mutual information of a state  $I(A : B)_{\rho} := D(\rho_{AB} || \rho_A \otimes \rho_B)$ .

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- $\odot I_{\max}^{\varepsilon}(A:B)_{\rho}$  is useful in quantum information tasks (state redistribution...)
- We will generalize these notations and results to a channel's version and find their connection with the quantum channel simulation task.

#### Definition

For any quantum channel  $\mathbb{N}_{A' \to B}$  we define the max-information of the channel  $\mathbb{N}$  as

$$I_{\max}(A:B)_{\mathcal{N}} := I_{\max}(A:B)_{\mathcal{N}_{A'\to B}(\Phi_{AA'})}, \qquad (14)$$

where  $\Phi_{AA'}$  is the maximally entangled state.

Note: We can replace  $\Phi_{AA'}$  to any pure state  $\phi_{AA'}$  with Schmidt rank |A'|.

#### Definition

The channel's smooth max-information is defined by

$$I_{\max}^{\varepsilon}(A:B)_{\mathcal{N}} := \inf_{\substack{\frac{1}{2} \|\widetilde{\mathcal{N}} - \mathcal{N}\|_{0} \leq \varepsilon \\ \widetilde{\mathcal{N}} \in CPTP(A':B)}} I_{\max}(A:B)_{\widetilde{\mathcal{N}}}, \qquad (15)$$

The channel's smooth max-information is **monotone** under composition with CPTP maps, i.e., for any CPTP maps  $\mathcal{N}_{A'_1 \to B_1}$ ,  $\mathcal{F}_{A'_0 \to A'_1}$  and  $\mathcal{T}_{B_1 \to B_0}$ ,

$$I_{\max}^{\varepsilon} \left( A_0 : B_0 \right)_{\mathfrak{T} \circ \mathcal{N} \circ \mathfrak{F}} \le I_{\max}^{\varepsilon} \left( A_1 : B_1 \right)_{\mathcal{N}}.$$
(16)

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#### Theorem

*For any quantum channel*  $\mathbb{N}_{A' \to B}$  *and given error tolerance*  $\varepsilon \ge 0$ *, we have* 

$$S_{\text{NS},\varepsilon}^{(1)}(\mathcal{N}) = \frac{1}{2} I_{\text{max}}^{\varepsilon} (A:B)_{\mathcal{N}}.$$
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The channel's smooth max-information has the asymptotic equipartition property,

$$\lim_{\varepsilon \to 0} \lim_{n \to \infty} \frac{1}{n} I^{\varepsilon}_{\max} \left( A : B \right)_{\mathcal{N}^{\otimes n}} = I \left( A : B \right)_{\mathcal{N}}, \tag{18}$$

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Sketch of proof:

$$\frac{1}{2} \lim_{\epsilon \to 0} \lim_{n \to \infty} \frac{1}{n} I_{\max}^{\epsilon} (A : B)_{\mathcal{N}^{\otimes n}} = \lim_{\epsilon \to 0} \lim_{n \to \infty} \frac{1}{n} S_{NS,\epsilon}^{(1)} (\mathcal{N}^{\otimes n})$$

$$= S_{NS} (\mathcal{N}) \qquad [by definition]$$

$$= Q_E (\mathcal{N}) \qquad [Bennett et al., 2014]$$

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Note: on the other hand, if we can proof Eq. (18) directly, it implies that

$$Q_{\rm E}(\mathcal{N}) = Q_{\rm NS}(\mathcal{N}) = S_{\rm NS}(\mathcal{N}). \tag{19}$$



Directly prove the AEP

$$\lim_{\varepsilon \to 0} \lim_{n \to \infty} \frac{1}{n} I_{\max}^{\varepsilon} (A:B)_{\mathcal{N}^{\otimes n}} = I(A:B)_{\mathcal{N}}. \quad \checkmark \text{ online soon}$$

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◎ Second-order asymptotics for 
$$C_{E}(\mathcal{N}) (= 2 \cdot Q_{E}(\mathcal{N}))$$
?

**Q:** What is the optimal rate  $r_{n,\varepsilon}$  to reliably transmit classical information via *n* uses of the quantum channel with entanglement assistance?

A second-order lower bound has been established [Datta et al., 2016]:

$$r_{n,\varepsilon} \ge C_{\mathrm{E}}(\mathbb{N}) + \sqrt{\frac{V_{\mathrm{E}}(\mathbb{N})}{n}} \Phi^{-1}(\varepsilon) + O\left(\frac{\log n}{n}\right).$$
(20)

It was conjectured that

$$r_{n,\varepsilon} = C_{\rm E}(\mathbb{N}) + \sqrt{\frac{V_{\rm E}(\mathbb{N})}{n}} \, \Phi^{-1}(\varepsilon) + o\left(\frac{\log n}{n}\right). \tag{21}$$

Obtaining the second-order asymptotics of  $I_{\max}^{\varepsilon}(A : B)_{\mathcal{N}}$  may provide a matching upper bound and solve the conjecture.

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• Other interesting applications of  $I_{\max}^{\varepsilon}(A:B)_{\mathcal{N}}$ ?

# Thanks for your attention!